

LETTER

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Mode-locked thulium fiber laser with MoS₂

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Abstract

Liquid-phase exfoliated 2D material multilayer MoS₂ is transferred onto a gold mirror and its saturable absorption at the 2 μm wavelength region is experimentally observed. This transferred MoS₂ saturable absorber has a modulation depth of 13.6% and a saturation intensity of 23.1 MW cm⁻². This saturable absorber is integrated into a linear Tm³⁺ fiber laser cavity, and stable fundamental-frequency mode-locking operation is realized at 2 μm with pulse energy of 15.5 nJ, pulse width of ~843 ps, and a repetition rate of 9.67 MHz. The laser spectral width is ~17.3 nm with a center wavelength of 1905 nm. This first presence of mode-locking with multilayer MoS₂ sheets in the 2 μm wavelength region verifies that multilayer MoS₂ is a good candidate for broadband mode-locking comparable to graphene, as well as a good mode-locker for achieving high pulse energy.

Keywords: thulium-doped fiber laser, mode-locking, multilayer MoS₂, dissipative soliton

(Some figures may appear in colour only in the online journal)

1. Introduction

Tm³⁺-doped fiber lasers (TDFLs) at 2 μm have received wide research interests in recent years due to their valuable applications in sensing, medical surgery, industrial machining and scientific experiments [1, 2]. To meet these application requirements, high-peak-power and/or high-energy 2 μm laser pulses are needed, which are generally produced by using Q-switching [3, 4] or mode-locking [5, 6] methods. Compared with Q-switching, mode-locking can provide much narrower pulse duration and higher peak power, as well as higher stability in the time domain.

Passive mode-locking, due to its simple configuration, has been a much preferred choice to obtain short pulses from 2 μm TDFLs, especially with the maturely developed semiconductor saturable absorbers (SAs) as modulators [7]. However, semiconductor SA has some drawbacks such as a complex design for improving damage threshold [8], a narrow working

wavelength range, etc. Recently, graphene (a monolayer of two-dimensional (2D) carbon atoms in a honeycomb structure) has attracted great attention and has been extensively used for mode-locking of TDFLs [9, 10] due to its advantages of large absorption per layer [11], wide operation spectral range [12], and moderate damage threshold and ultrafast recovery time [13]. Another kind of 2D material MoS₂ has also attracted wide attention in the fiber laser community, and has been extensively explored to mode lock fiber lasers [14–18]. Owing to the semiconductor characteristics of layered MoS₂, especially the monolayer MoS₂ a direct band semiconductor, the band gap determines the energy of photons to be absorbed. However, studies have proven that layered MoS₂, in particular when stoichiometric defects (non-ideal atomic ratio) are introduced, also possesses wideband absorption and saturable absorption features [19], which are prerequisites for mode-locking of lasers. Thereafter, extensive researches have been dedicated to exploring of mode-locking operation and

related characteristics of fiber lasers in the $1\ \mu\text{m}$ [14, 15] and $1.5\ \mu\text{m}$ [16–18] wavelength regions. However, mode-locking with MoS_2 in the $2\ \mu\text{m}$ regime has not been reported to date. Owing to large anomalous dispersion of the gain fiber and lower absorption of layered MoS_2 at $2\ \mu\text{m}$, mode-locking operation with this kind of 2D material and corresponding behaviors in the near mid-infrared spectral region still need further verification and investigation.

In this letter, we report $2\ \mu\text{m}$ Tm^{3+} fiber lasers mode-locked with multilayer MoS_2 with 15.5 nJ pulse energy. Simple linear cavity incorporated with gain fiber and the multilayer MoS_2 modulator is constructed to achieve mode-locking operation, showing that mode-locking capability of layered MoS_2 sheets can be definitely extended to the $2\ \mu\text{m}$ wavelength region. With normal net cavity dispersion, fundamental mode-locking in the dissipative soliton regime is achieved with maximum output power of $>150\ \text{mW}$, pulse width of 843 ps and pulse repetition rate of 9.67 MHz. This is the first time mode-locking in the $2\ \mu\text{m}$ wavelength region with 2D MoS_2 materials has been realized.

2. Material characterization and experimental setup

The multilayer MoS_2 was synthesized with the liquid-phase exfoliation method (LPE). The MoS_2 powder was dispersed in ethanol solvent and sonicated for 1 h to produce few-layer MoS_2 nano-sheets. The dispersions were then put for 5 h so that the unexfoliated large-size flakes can sink to the bottom and be separated. The top 2/3 dispersions were collected for experimental use and separately transferred onto one side of a gold mirror by drop coating and air drying. After transfer, the MoS_2 on the mirror was detected by Raman spectroscopy (Renishaw Invia, 488 nm) and saturable absorption test. A representative Raman spectrum of the transferred sample is shown in figure 1(a). E_{2g}^1 and A_{1g} modes are nearly coincident with those of bulk MoS_2 ($\sim 383\ \text{cm}^{-1}$ for E_{2g}^1 and $\sim 408\ \text{cm}^{-1}$ for A_{1g}), showing that the MoS_2 sample has a thickness of ~ 4 layers [20]. The transmittance of the multilayer MoS_2 on the gold mirror is shown in figure 1(b). Here, we used a reflection method to measure the saturable absorption of the sample, i.e. the incident probe laser beam was incident on the sample, transited the MoS_2 sheets, reflected by the gold mirror, then transited the MoS_2 sheets once more, and finally the retro-reflected laser beam was measured. A self-constructed ps ($\sim 800\ \text{ps}$) fiber laser at 1940 nm with a repetition rate of $\sim 55\ \text{MHz}$ was used as the input pulse source. The power-dependent absorption measurement was carried out through changing the input laser power. To obtain the nonlinear optical parameters, a simple saturable absorption model of $T(I) = 1 - \alpha_0 \times \exp(-I/I_{\text{sat}}) - \alpha_{\text{ns}}$ [14] was used to fit the measured data. $T(I)$ is the transmission, α_0 is the modulation depth, I is the input intensity, I_{sat} is the saturation intensity, and α_{ns} is the non-saturable absorbance. As shown in figure 1(b), the modulation depth α_0 , non-saturable loss α_{ns} and saturation intensity I_{sat} were obtained to be 13.6%, 16.7% and $23.1\ \text{MW cm}^{-2}$, respectively. The modulation depth achieved here (at the $2\ \mu\text{m}$

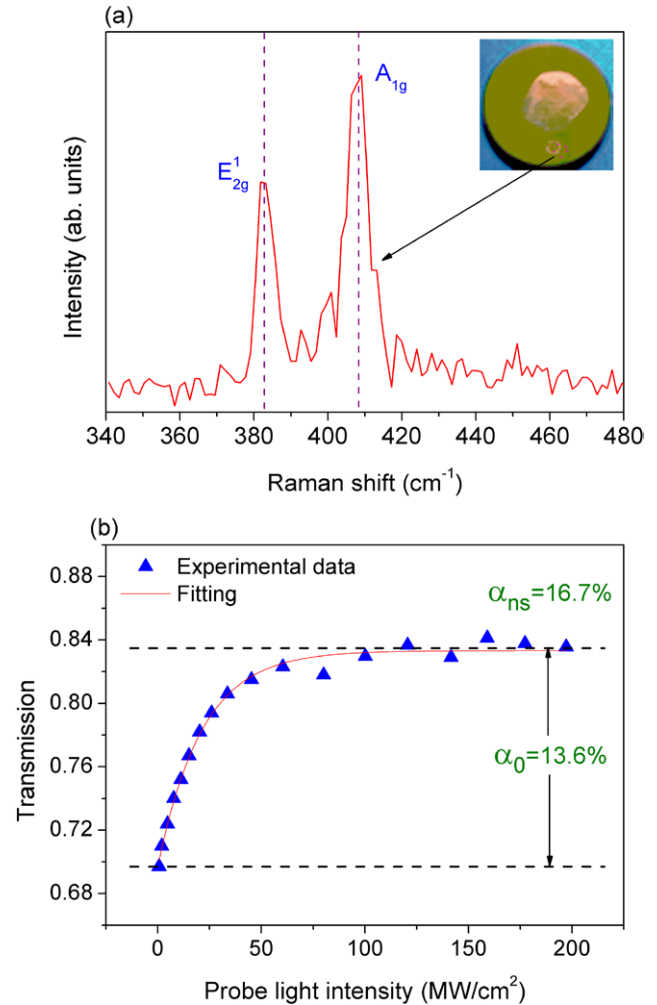


Figure 1. Raman spectrum of the adopted multilayer MoS_2 sheets and (b) nonlinear absorption of the multilayer MoS_2 sheets coated on a gold mirror.

wavelength region) is comparable to that measured in the $1\ \mu\text{m}$ [14, 15] region and much larger than that in the $1.5\ \mu\text{m}$ [17, 18] wavelength region. This large modulation depth of the MoS_2 saturable absorber at the $2\ \mu\text{m}$ wavelength region can be made advantage of to suppress wave breaking in mode-locking operation [16, 21], and thus indicates that this sample can be used as an effective saturable absorber in the $2\ \mu\text{m}$ wavelength range to realize large pulse energy.

The experimental setup for the MoS_2 mode-locked TDFL is shown in figure 2. The gain fiber was 12 cm highly doped single-cladding Tm^{3+} -doped silica fiber ($5/125\ \mu\text{m}$, 0.24 NA) with core absorption of $\sim 350\ \text{dB m}^{-1}$ at $\sim 1550\ \text{nm}$. The pump source was a CW erbium/ytterbium-codoped fiber laser (EYFL) with maximum output of 1 W centered at 1550 nm. The pump light was delivered into the gain fiber by a 1550/2000 nm wavelength division multiplexing coupler (WDM), which is made by SMF-28 fiber. After the gain fiber, a piece of 4.6 m dispersion compensating fiber (DCF) ($2.2\ \mu\text{m}$, 0.35 NA core) was used to provide total net normal cavity dispersion for achieving dissipative soliton mode-locking. At the output end, a piece of 4 m SMF-28 fiber was spliced to decrease the fundamental mode-locking frequency.

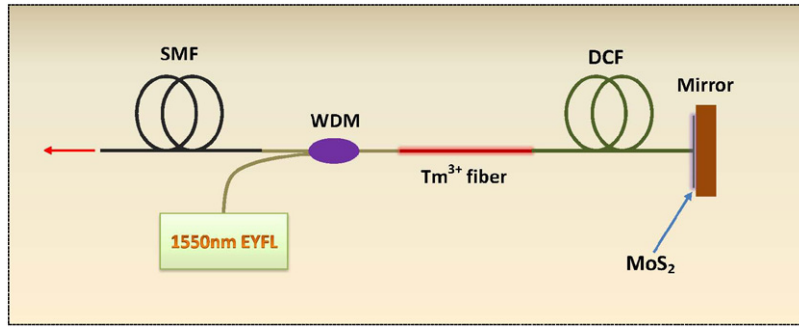


Figure 2. Experimental setup of the mode-locked Tm^{3+} fiber laser. LD: laser diode; SMF: single-mode fiber; DCF: dispersion compensating fiber; EYFL: erbium/ytterbium-codoped fiber laser; WDM: wavelength division multiplexing coupler.

The dispersions of the DCF fiber, the SMF-28 fiber (including 1.62 m pigtail SMF-28 fiber) and the Tm^{3+} fiber at $1.9\ \mu\text{m}$ are $93, -67$ and $-12\ \text{ps}^2\text{km}^{-1}$, respectively [22], which give a total net normal cavity dispersion of $\sim 0.05\ \text{ps}^2$. At the rear end, a high-reflection gold mirror (coated with the MoS_2 sheets) provided feedback, and this mirror combined with the perpendicularly cleaved output-end fiber facet ($\sim 3.5\%$ Fresnel reflection) completed the laser cavity. The gold mirror, on which the MoS_2 was adsorbed, was fixed on a five-dimension stage (3D translation and 2D rotation) for fine position adjustment and the fiber end was perpendicularly butt coupled to the MoS_2 sheets.

Laser output power was measured with a power meter (FieldMate, Coherent Co.) and the laser spectrum was tested with a mid-infrared analyzer (SIR 5000, SandHouse Co.) with a spectral resolution of $0.22\ \text{nm}$. Laser pulsing dynamics were recorded with a $2.5\ \text{GHz}$ Agilent oscilloscope (DSO9254A) combined with a $3.5\ \text{GHz}$ InGaAs detector. The radio-frequency spectrum was measured with a $26.5\ \text{GHz}$ microwave spectrum analyzer (N9938A, FieldFox).

3. Results and discussions

When the 2 micron laser operation was initiated with pump power over threshold ($\sim 430\ \text{mW}$), the laser first went to continuous wave (CW) operation. When the pump power was increased to over $630\ \text{mW}$ and the MoS_2 sheet position was carefully adjusted, the laser came to the stable Q-switching regime. Then, further raising the pump power to over $700\ \text{mW}$, stable mode-locking operation of the TDFL occurred, which could be sustained up to the available maximum pump power ($730\ \text{mW}$). The output power shows a linear increase with pump power, and the maximum output power is $150\ \text{mW}$, as shown in figure 3. The slope efficiency is 43.6% with respect to pump power. The laser spectrum of the mode-locked TDFL at the maximum output level is present in the inset of figure 3. The spectrum is centered at $\sim 1905\ \text{nm}$ with an FWHM (full width at half maximum) bandwidth of about $17.3\ \text{nm}$. This spectral width is much larger than that of the 1 and $1.5\ \mu\text{m}$ counterparts [14–17], showing that much narrower mode-locked pulse duration can be feasible if dispersion is appropriately compensated.

The laser pulse trains obtained at the maximum output levels are shown in figure 4(a). The $103.4\ \text{ns}$ period time corresponds well to the cavity round trip time (the total fiber

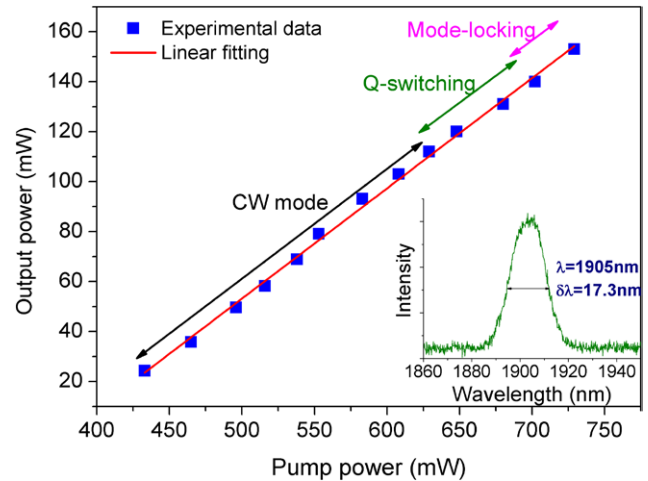


Figure 3. Output of the mode-locked Tm^{3+} fiber laser. The laser threshold is around $430\ \text{mW}$, and the slope efficiency is about 43.6% . Square dots are measured data and the solid line is linear fitting. Inset shows the laser spectrum of the mode-locked Tm^{3+} fiber laser (in linear scale).

length is $\sim 10\ \text{m}$), showing that the mode-locking operates at the fundamental frequency of $9.67\ \text{MHz}$. The intensity stability between different pulses is $>95\%$. Considering the output power of $150\ \text{mW}$, the output laser pulse has single pulse energy of $15.5\ \text{nJ}$. This is the highest pulse energy ever achieved in mode-locked fiber lasers with MoS_2 modulators, and this also demonstrates that 2D material MoS_2 has great potential in high power photoelectronics and integrated photonics.

Single pulse characteristics measured at the maximum power level is shown in figure 4(b). The single pulse shows a nearly Gaussian shape with FWHM pulse width of about $843\ \text{ps}$. This pulse width is comparable to that obtained from the $1\ \mu\text{m}$ dissipative-soliton counterparts [14, 15]. This pulse duration, combined with the above mentioned spectral width, gives a time-bandwidth product of ~ 1200 , which is by far larger than the Fourier transform limited one of 0.441 if a Gaussian pulse shape is assumed. This means that the mode-locked laser pulse is highly chirped, which also accounts for the experimentally observed high pulse energy. This large mode-locked pulse duration originates from the unique characteristics of dissipative solitons, chirping pulse to achieve high pulse energy. In fact, 2D MoS_2 (monolayer or few-layer) has ultrafast recovery times of tens of femtoseconds [23] and

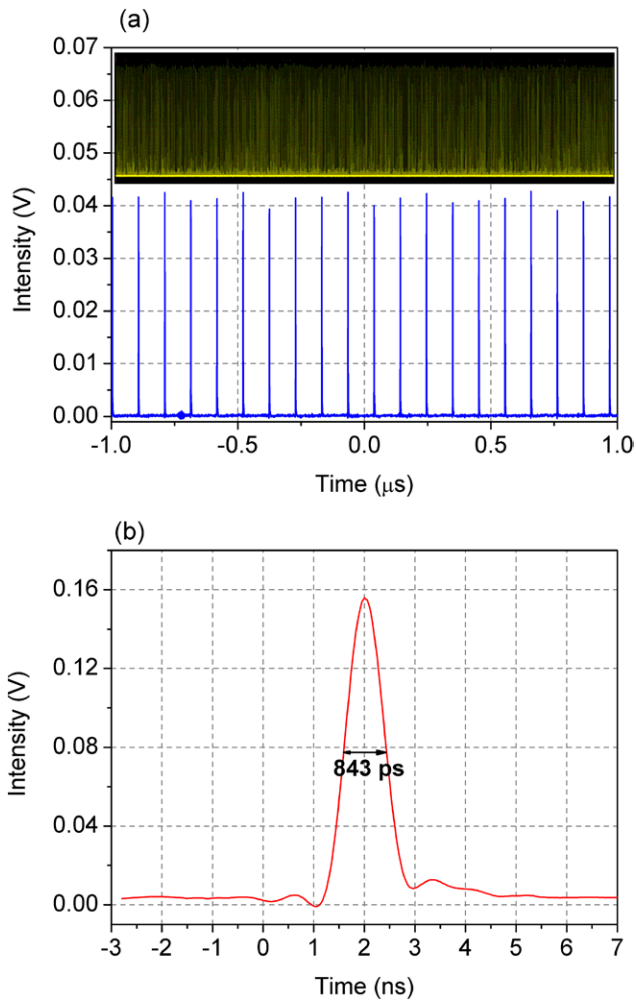


Figure 4. Laser pulse train (a) and single pulse (b) of the MoS₂ mode-locked Tm³⁺ fiber laser. Insets show the oscilloscope traces.

~100 ps [24, 25], corresponding respectively to the intraband transition and interband transition of excited free carriers. Recently, mode-locking with similar multilayer MoS₂ saturable absorbers has achieved femtosecond time-scale pulse durations [18, 26, 27]. Based on the pulse spectral width (17.3 nm) in our experiment, a Fourier transform-limited pulse width of ~247 fs is expected provided that the entire pulse chirp can be compensated. To out-cavity compensate the high pulse chirp of our solitons (843 ps), coarse calculation shows that around 1.1 km single-mode fiber (SMF-28 with dispersion of $-67 \text{ ps}^2 \text{ km}^{-1}$) is needed.

In order to show the stability of the mode-locking state, the radio-frequency spectrum was measured with a bandwidth resolution of 0.1 MHz and the results are shown in figure 5. Over the 100 MHz range, only the fundamental frequency of 9.67 MHz is present, and the signal-to-noise ratio is >55 dB, showing high stability of the mode-locking operation with the as manufactured MoS₂ multilayer sheets.

4. Summary

In conclusion, stable mode-locked 2 μm TDFL with multilayer MoS₂ sheets has been realized with 150 mW output power in

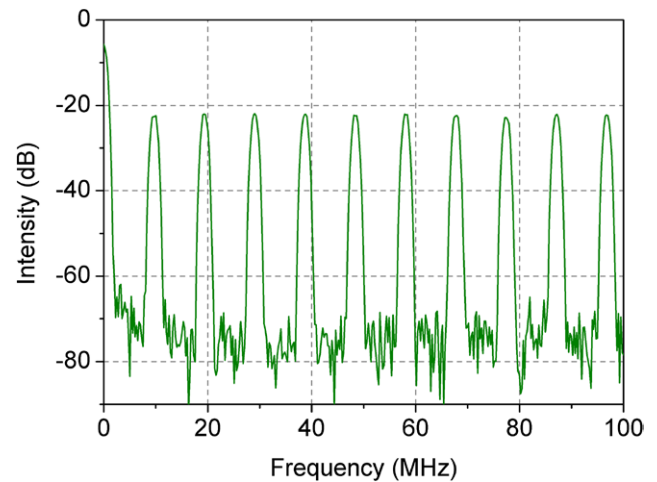


Figure 5. Radiofrequency spectral profile of the mode-locked Tm³⁺ fiber laser.

a compact linear cavity. Fundamental-frequency mode-locking operation provides >15 nJ pulse energy and ~10 MHz repetition rate. This highly chirped dissipative soliton pulse has a temporal width of 843 ps, which can be dechirped outside the cavity. This first stable mode-locking operation with multilayer MoS₂ in the 2 μm region clearly indicates that few-layer transition metal dichalcogenides can also function as wide-band optoelectronic elements, comparable to graphene. The achieved >15 nJ pulse energy, in addition, shows that this 2D material can be high power photonic devices, even in the 2 μm wavelength region.

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